# Nano-engineered Casimir forces

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## Collaborators



## Theory:

Engineering the Casimir force with metamaterials: Peter Milonni (LANL)

Felipe da Rosa (LANL)

Engineering the Casimir force with geometry: Paulo Maia Neto (Rio)

Serge Reynaud (Paris)

## Experiment:

Metamaterials for Casimir force: Toni Taylor (LANL)

Hou-Tong Chen (LANL)

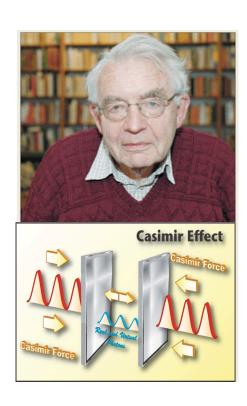
Casimir force measurements: Steve Lamoreaux (Yale)

Ricardo Decca (Indiana)

Roberto Onofrio (Dartmouth)

## The Casimir force





Casimir forces originate from changes in quantum vacuum fluctuations imposed by surface boundaries

They were predicted by the Dutch physicist Hendrik Casimir in 1948

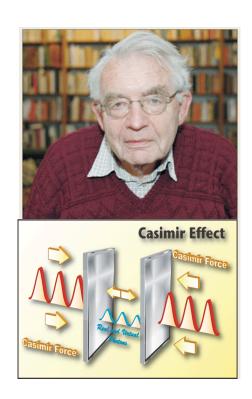
Dominant interaction in the micron and sub-micron lengthscales

$$\frac{F}{A} = \frac{\pi^2}{240} \; \frac{\hbar c}{d^4}$$

 $(130 \text{nN/cm}^2 @ d = 1 \mu\text{m})$ 

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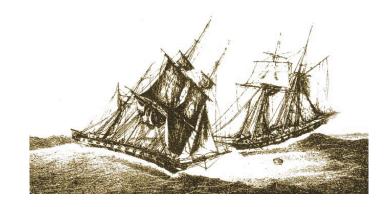
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$$\frac{F}{A} = \frac{\pi^2}{240} \; \frac{\hbar c}{d^4}$$

$$(130 \text{nN/cm}^2 @ d = 1 \mu \text{m})$$

Classical Analog: L'Album du Marin (1836)



# Relevant applications



Gravitation / Particle theory:

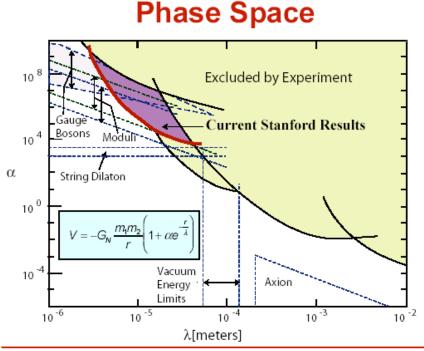
Some theories of particle physics predict deviations from the Newtonian gravitational potentials in the micron and submicron range

The Casimir force is the main background force to measure these

non-Newtonian corrections to gravity

Yukawa-like potential:

$$V(r) = -G\frac{m_1 m_2}{r} \left( 1 + \alpha \ e^{-r/\lambda} \right)$$



# Relevant applications



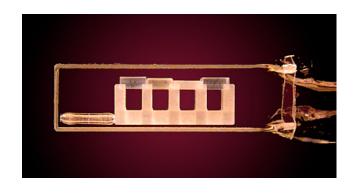
Quantum Science and Technology:

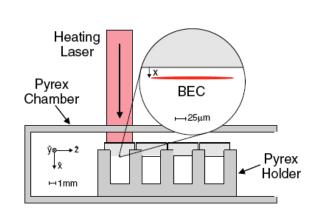
Atom-surface interactions

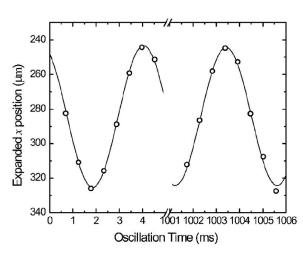
Precision measurements

Example: Casimir-Polder interaction between a BEC and a surface

E. Cornell et al, Phys. Rev. Lett. 98, 063201 (2007)







$$\gamma_x \equiv \frac{\omega_x - \omega_x'}{\omega_x} \simeq -\frac{1}{2\omega_x^2 m} \partial_x^2 U^*$$

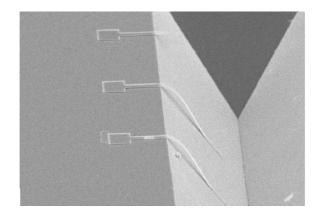
# Relevant applications



## Nanotechnology:

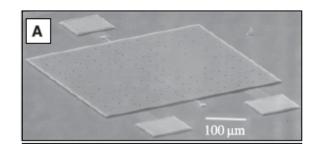
Problems with stiction of movable parts in MEMS

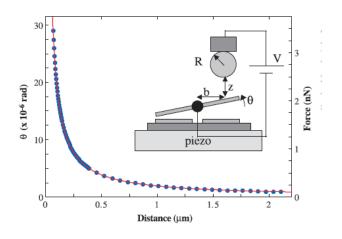
"pull-in" effect



Zhao et al, Adhesion Sci. Technol. 17, 519 (2003)

# Actuation in NEMS and MEMS driven by Casimir forces



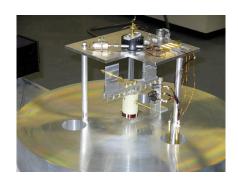


F. Capasso et al, Science 291, 1941 (2001)



## Torsion pendulum



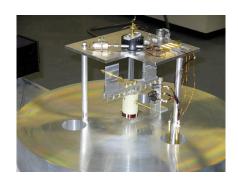


sphere-plane, d=1-10 um Lamoreaux



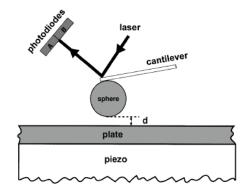
## Torsion pendulum





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## Atomic force microscope

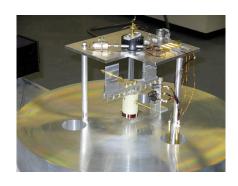


sphere-plane, d=200-1000 nm Mohideen



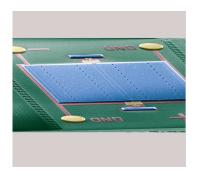
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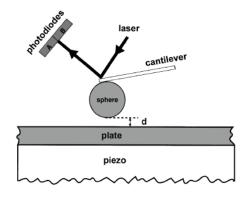
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### MEMS and NEMS



sphere-plane, d=200-1000 nm Capasso, Decca

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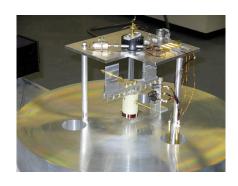


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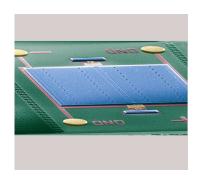
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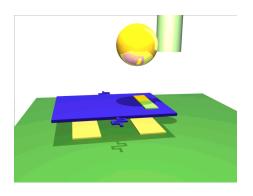




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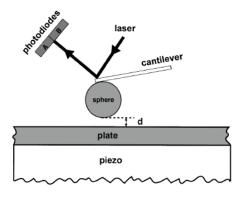
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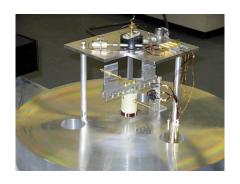


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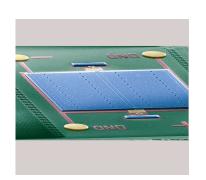
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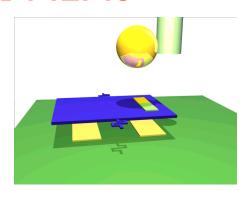




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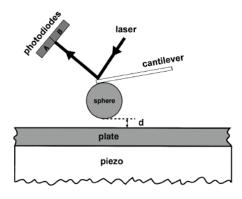
#### MEMS and NEMS





sphere-plane, d=200-1000 nm Capasso, Decca

## Atomic force microscope



sphere-plane, d=200-1000 nm Mohideen

#### Micro-cantilever



plane-plane, cylinder-plane, d=1-3 um Onofrio

# Tailoring the Casimir force



The magnitude and sign of the Casimir force depend on the geometry and composition of surfaces

Engineer geometries and designer materials for various applications:

- Demonstration of strongly modified/repulsive Casimir forces
- Demonstration of vacuum drag via lateral Casimir forces
- From ideal to real materials: The Lifshitz formula

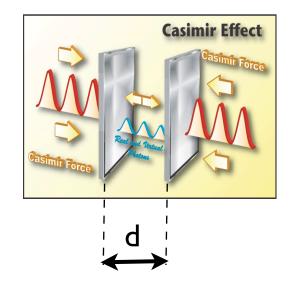
$$\frac{F}{A} = \frac{\pi^2}{240} \frac{\hbar c}{d^4} \qquad \longrightarrow \qquad \frac{F}{A} = 2k_B T \sum_{n=0}^{\infty} \int_{\xi_n/c}^{\infty} \frac{d\kappa}{2\pi} \kappa^2 \sum_{\lambda = \text{TE,TM}} \left( \frac{e^{2\kappa d}}{r_{\lambda_1} r_{\lambda_2}} - 1 \right)^{-1}$$

Dominant frequencies in the near-infrared/optical region of the EM spectrum (gaps d= 200-1000 nm)

## The Lifshitz formula



$$\frac{F}{A} = 2k_B T \sum_{n=0}^{\infty} \int_{\xi_n/c}^{\infty} \frac{d\kappa}{2\pi} \kappa^2 \sum_{\lambda = \text{TE,TM}} \left( \frac{e^{2\kappa d}}{r_{\lambda_1} r_{\lambda_2}} - 1 \right)^{-1}$$



$$\omega_n = i\xi_n = 2\pi i n k_B T/\hbar$$
 Matsubara frequencies

$$r_{\text{TM}} = \frac{\epsilon(i\xi_n)c\kappa - \sqrt{\xi_n^2[\epsilon(i\xi_n)\mu(i\xi_n) - 1] + \kappa^2 c^2}}{\epsilon(i\xi_n)c\kappa + \sqrt{\xi_n^2[\epsilon(i\xi_n)\mu(i\xi_n) - 1] + \kappa^2 c^2}}$$

Reflection coefficients

$$r_{\rm TE} = r_{\rm TM} \text{ with } \epsilon \leftrightarrow \mu$$

#### Kramers-Kronig (causality) relations:

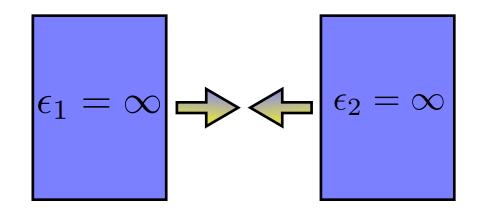
$$\epsilon(i\xi_n) = 1 + \frac{2}{\pi} \int_0^\infty \frac{\omega \epsilon''(\omega)}{\omega^2 + \xi_n^2} d\omega \qquad \qquad \mu(i\xi_n) = 1 + \frac{2}{\pi} \int_0^\infty \frac{\omega \mu''(\omega)}{\omega^2 + \xi_n^2} d\omega$$



#### Ideal attractive limit

Casimir 1948

$$\frac{F}{A} = +\frac{\pi^2}{240} \frac{\hbar c}{d^4}$$

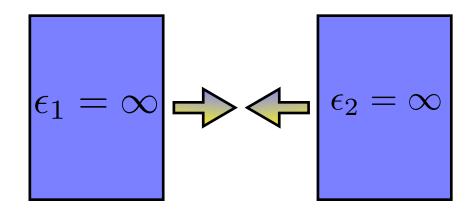




Ideal attractive limit

Casimir 1948

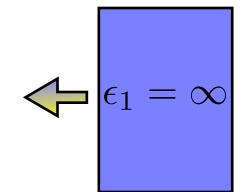
$$\frac{F}{A} = +\frac{\pi^2}{240} \frac{\hbar c}{d^4}$$

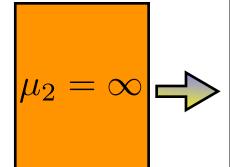


Ideal repulsive limit

Boyer 1974

$$\frac{F}{A} = -\frac{7}{8} \frac{\pi^2}{240} \frac{\hbar c}{d^4}$$



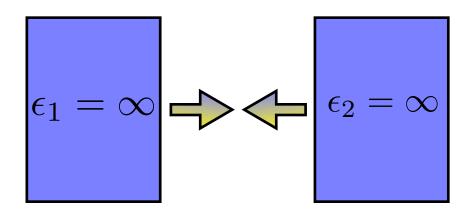




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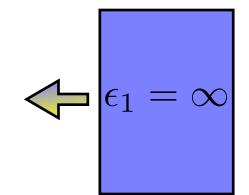
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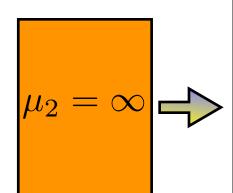


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$$\frac{F}{A} = -\frac{7}{8} \frac{\pi^2}{240} \frac{\hbar c}{d^4}$$





Real repulsive limit

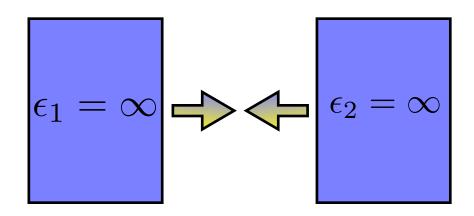
Casimir repulsion is associated with strong electric-magnetic interactions. However, natural occurring materials do NOT have strong magnetic response in the optical region, i.e.  $\mu=1$ 



Ideal attractive limit

Casimir 1948

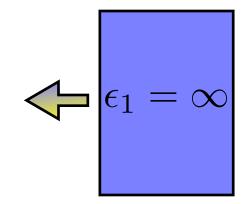
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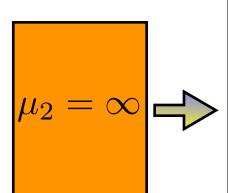


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Real repulsive limit

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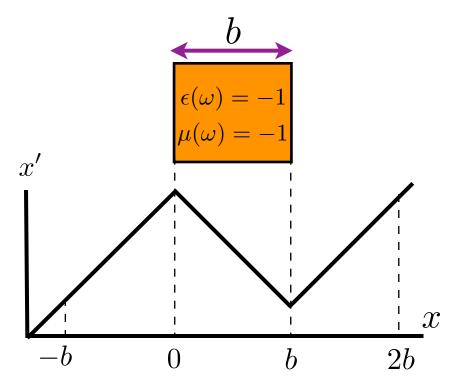
Physicists have 'solved' mystery of levitation - Telegraph

http://www.telegraph.co.uk/news/main.jhtml?xml=/news/2007/08/0...





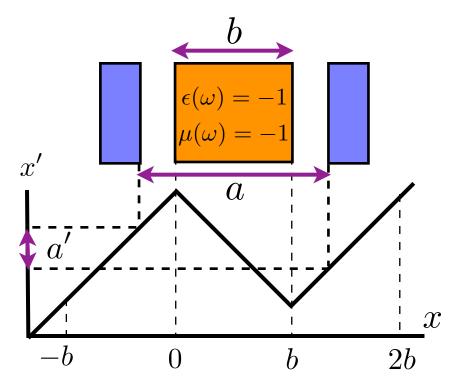
Transformation media Leonhardt et al, NJP 9, 254 (2007)



Perfect lens: EM field in -b < x < 0 is mapped into x'. There are two images, one inside the device and one in b < x < 2b.



Transformation media Leonhardt et al, NJP 9, 254 (2007)



Perfect lens: EM field in -b<x<0 is mapped into x'. There are two images, one inside the device and one in b<x<2b.

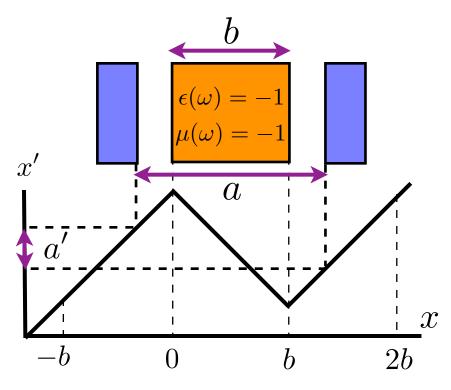
Casimir cavity: 
$$a' = |a - 2b|$$

When a < 2b (plates within the imaging range of the perfect lens)

$$\Rightarrow f = -\frac{\partial U}{\partial a'} \frac{\partial a'}{\partial a} = +\frac{\hbar c \pi^2}{240 a'^4} \Rightarrow \text{Repulsion}$$



### Transformation media Leonhardt et al, NJP 9, 254 (2007)



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Casimir cavity: 
$$a' = |a - 2b|$$

When a < 2b (plates within the imaging range of the perfect lens)

$$\frac{1}{2b}^{x}$$
  $\implies$   $f = -\frac{\partial U}{\partial a'} \frac{\partial a'}{\partial a} = +\frac{\hbar c \pi^2}{240a'^4} \implies$  Repulsion

For real materials, however .....

- According to causality, no passive medium (  $\epsilon$ " ( $\omega$ ) > 0 ) can sustain  $\epsilon$ ,  $\mu \simeq -1$  over a wide range of frequencies. In fact,  $\epsilon(i\xi), \mu(i\xi) > 0$
- Leonhardt proposes to use an active MM (  $\epsilon$ " ( $\omega$ ) < 0 ) in order to get repulsion. But then the whole approach breaks down, as real photons would be emitted into the quantum vacuum.



#### Drude-Lorentz model:

$$\epsilon_{\alpha}(\omega) = 1 - \frac{\Omega_{E,\alpha}^2}{\omega^2 - \omega_{E,\alpha}^2 + i\Gamma_{E,\alpha}\omega}$$
$$\mu_{\alpha}(\omega) = 1 - \frac{\Omega_{M,\alpha}^2}{\omega^2 - \omega_{M,\alpha}^2 + i\Gamma_{M,\alpha}\omega}$$

#### Typical separations

$$d = 200 - 1000 \text{ nm}$$

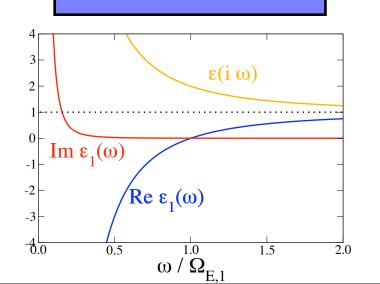


#### Infrared-optical frequencies

$$\Omega/2\pi = 5 \times 10^{14} \, \mathrm{rad \, s^{-1}}$$

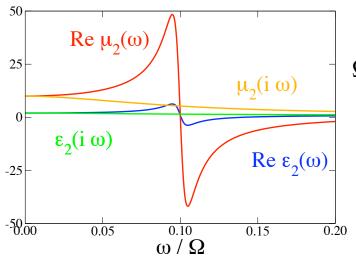
#### Drude metal (Au)

$$\Omega_E = 9.0 \; \mathrm{eV} \; \; \Gamma_E = 35 \; \mathrm{meV}$$



#### Metamaterial

Re 
$$\epsilon_2(\omega) < 0$$
 Re  $\mu_2(\omega) < 0$ 



$$\Omega_{E,2}/\Omega = 0.1$$
  $\Omega_{M,2}/\Omega = 0.3$ 

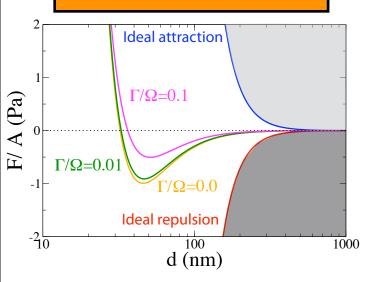
$$\omega_{E,2}/\Omega = \omega_{M,2}/\Omega = 0.1$$

$$\Gamma_{E,2}/\Omega = \Gamma_{M,2}/\Omega = 0.01$$



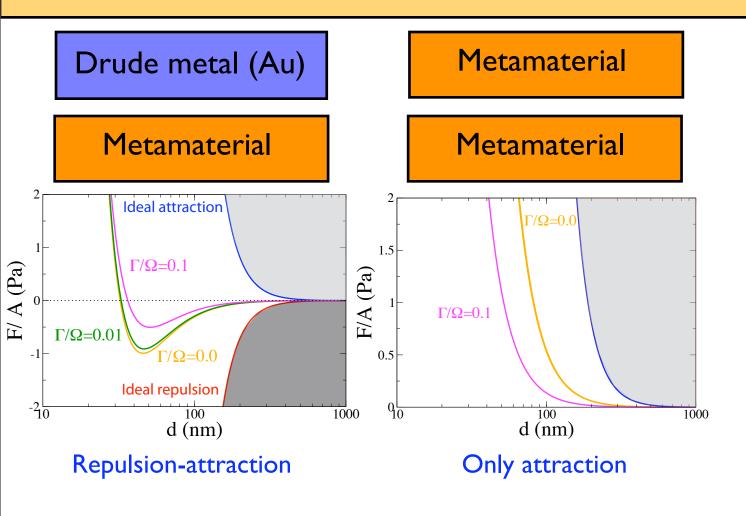
Drude metal (Au)

#### Metamaterial

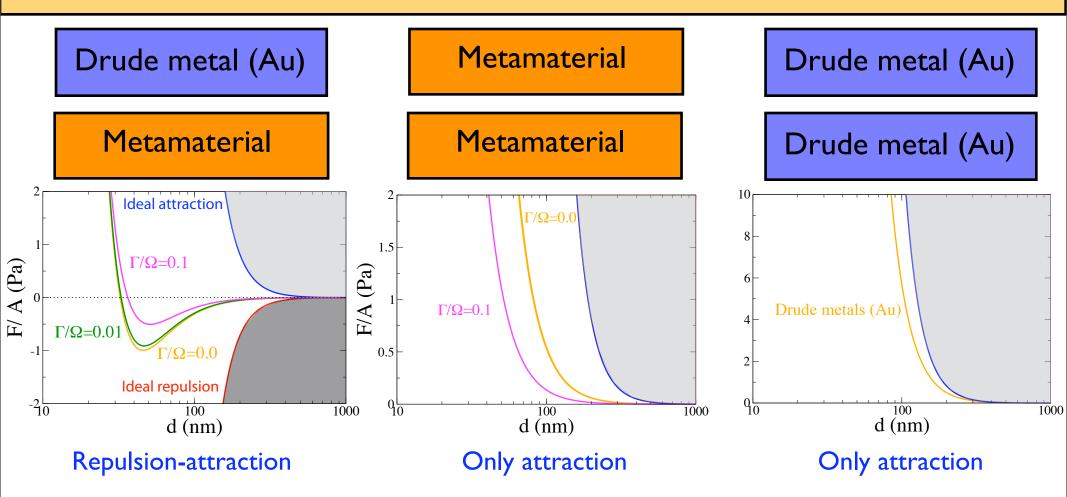


Repulsion-attraction

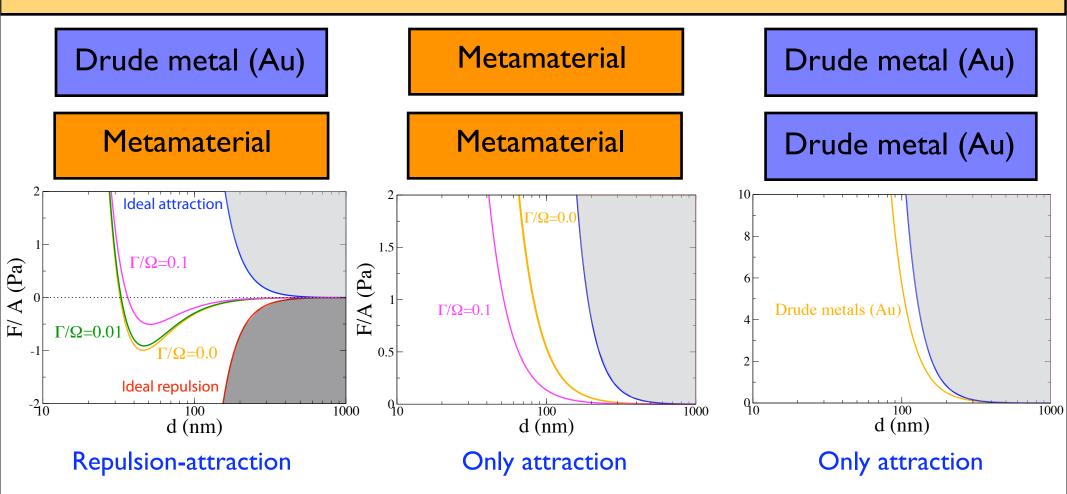












A slab made of Au ( $\rho=19.3~{\rm gr/cm^3}$ ) of width  $\delta=1\mu{\rm m}~{\rm could}$  levitate in front of one of these MMs at a distance of  $d\approx110~{\rm nm}~!!!$ 

Casimir and metamaterials Henkel et al, EPL 72, 929 (2005)
Casimir and surface plasmons Lambrecht et al, PRL 94, 110404 (2005)
van der Waals in magneto-dielectrics Dalvit et al, PRA 75, 052117 (2007)

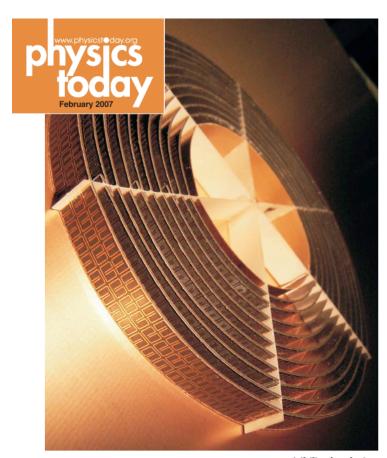
## Conclusions



- Metamaterials can strongly influence the quantum vacuum, providing a route towards quantum levitation.
- ☐ Build MMs with strong magnetic response at infrared-optical frequencies, corresponding to gaps between 200 nm and 10 microns.
- ☐ Validity of the continuum description? Extensions of Lifshitz formula to periodic structures...
- ☐ Fabrication of MMs on torsion pendulum plane or on oscillating MEMS/NEMS
- ☐ Similar issues apply to tailored plasmons, e.g. SAA

## The Casimir force and MMs





Invisibility by design

Review article by Steve Lamoreaux on page 40:

"Casimir forces: still surprising after 60 years"